Hunting Natural SUSY at LHC

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EW Fine-tuning

arXiv: 1207.3343, 1404.1386

Minimization condition for higgs scalar potential (1-loop)

$$\frac{M_Z^2}{2} = \frac{(m_{H_d}^2 + \Sigma_d^d) - (m_{H_u}^2 + \Sigma_u^u) \tan^2 \beta}{\tan^2 \beta - 1} - \mu^2$$

• Sparticles contribute via radiative corrections inside Σ

$$\Sigma_u^u(\tilde{t}_{1,2}) = \frac{3}{16\pi^2} F(m_{\tilde{t}_{1,2}}^2) \times \left[f_t^2 - g_Z^2 \mp \frac{f_t^2 A_t^2 - 8g_Z^2 (\frac{1}{4} - \frac{2}{3}x_W) \Delta_t}{m_{\tilde{t}_2}^2 - m_{\tilde{t}_1}^2} \right]$$
$$F(m^2) = m^2 \left(\log(m^2/Q^2) - 1 \right), \text{ with } Q^2 = m_{\tilde{t}_1} m_{\tilde{t}_2}$$

Naturalness requires no large cancellations ⇒ fine-tuning parameter

$$\Delta_{EW} \equiv \max\left(\frac{m_{H_u}^2}{M_Z^2/2} \frac{\tan^2 \beta}{\tan^2 \beta - 1}, \frac{\Sigma_u^u}{M_Z^2/2} \frac{\tan^2 \beta}{\tan^2 \beta - 1}, \frac{\mu^2}{M_Z^2/2}, \ldots\right)$$

- Limited value of Δ_{EW} \Rightarrow upper limit on $m_{H_u}^2$ and μ^2
- Δ_{EW} is determined by the spectrum and does not depend on high-scale dynamics

Properties of Δ_{EW}

arXiv: 1207.3343,

1404.1386

FT is a property of the model and can be quantified by measure

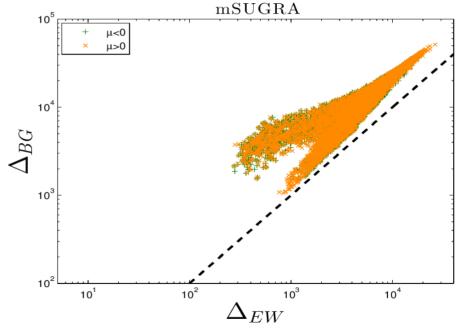
Ellis, Enqvist, Nanopoulos, Zwirner '86 $\Delta_{BG} \equiv max_i \left| \frac{\Delta M_Z^2/M_Z^2}{\Delta a_i/a_i} \right| = max_i \left| \frac{\partial \ln M_Z^2}{\partial \ln a_i} \right|$ Barbieri, Giudice '88

• Δ_{EW} provides a bound on FT, that can be saturated with specific

$$\Delta_{BG} > \Delta_{EW}$$

model para correlations

• If Δ_{EW} is large, any underlying theory that leads to the spectrum will be fine-tuned.

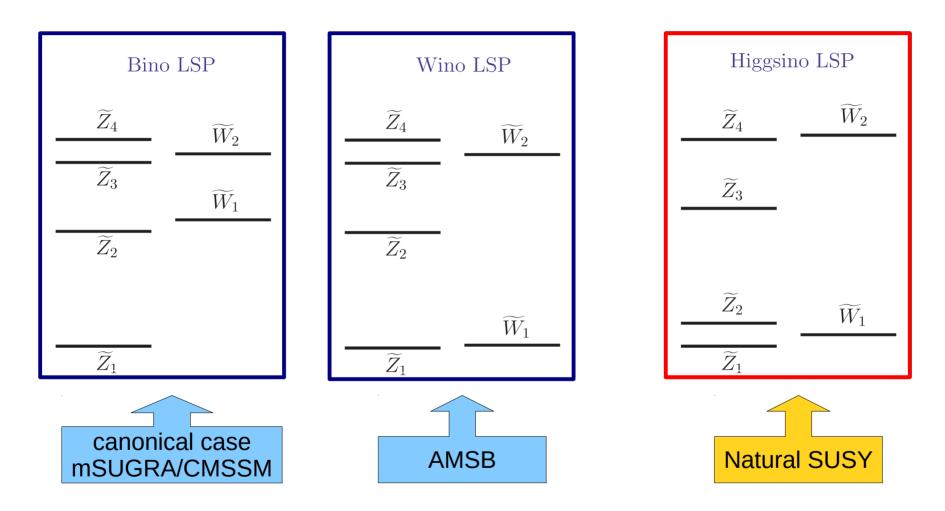


• Small Δ_{EW} does not imply no FT, but leaves open the possibility of finding a theory with low FT.

Summary of Natural spectrum

- For m_h ~ 125 GeV and $\Delta_{EW} < 30$:
 - μ ~ 100-300 GeV
 - stop_1 ~ 1-2 TeV, stop_2 = sbottom_1 ~ 2-4 TeV, highly mixed by large A_t
 - gluino ~ 1-5 TeV
 - 1st/2nd generation squarks ~ 1-10 TeV
 - sleptons ~ 1-30 TeV
- This can be realized in a simple extension of mSUGRA, NUHM2 $m_0,\ m_{1/2},\ A_0,\ \tan\beta,\ \mu,\ m_A$
- Here small $m_{H_u}^2 \simeq -M_Z^2$ and lighter stops are generated by RGE evolution, hence Radiatively-driven Natural SUSY (RNS)

EW-ino spectrum



 Current ATLAS/CMS limits done in simplified models and not applicable to the light higgsino case

Hard trilepton signals

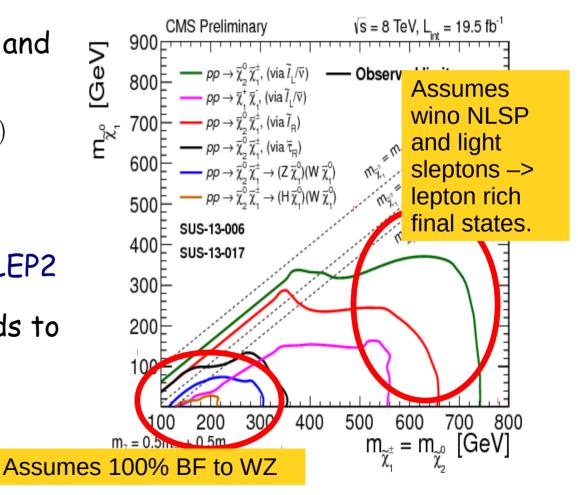
arXiv: 1310.4858

- Trileptons are golden channel for gaugino searches that relies on 3-body decay of neutralinos:
- In RNS sleptons are heavy and 3l only appear from

$$pp \to \widetilde{W}_2 \widetilde{Z}_4 \to (\widetilde{W}_1 Z) + (\widetilde{W}_1 W)$$



- No reach at LHC8 beyond LEP2
- At LHC14 the reach extends to $m_{1/2} = 500~(630)~{
 m GeV}$ for 300 (1000) fb⁻¹

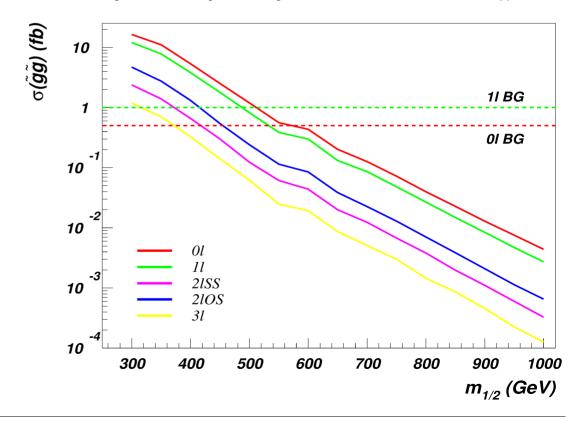


EW-inos in gluino cascades

arXiv: 1310.4858

- $pp o \tilde{g} \tilde{g} X$ followed by gluino cascade decays lead to multi-jets + multi-leptons + MET signature
- LHC14 can only probe part of natural gluino mass range, up to 1700 GeV (1900 GeV) for 300 (1000)fb⁻¹
- $\widetilde{Z}_2 \widetilde{Z}_1$ mass gap (if >25 GeV) can be measured from OS/SF dileptons

NUHM2: m_o =5 TeV, A_o =-1.6 m_o , $tan\beta$ =15, μ =150 GeV, m_A =1 TeV



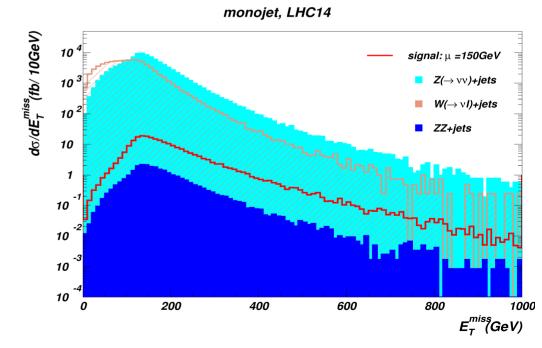
Mono-jets and mono-photons

arXiv: 1401.1162

- Higgsinos have compressed spectrum with mass gap 10-30 GeV only soft visible energy ==> higgsinos mostly appear as MET.
- $pp \to \widetilde{Z}_{1,2}\widetilde{Z}_{1,2}/\widetilde{Z}_{1,2}\widetilde{W}_1/\widetilde{W}_1\widetilde{W}_1 + (j \text{ or } \gamma)$
- Contact interaction, used in ATLAS/CMS search, not applicable: mediator mass $M_Z \ll \sqrt{\hat{s}} \sim p_T(jet) + E_T^{miss}$

leading to extra 1/s suppression for ME

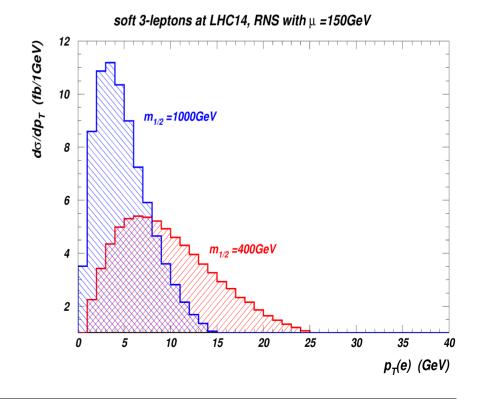
- Signal has same shape as BG and S/B ≈1%.
 - Detection is very challenging.



Soft trileptons

arXiv: 1302.5816

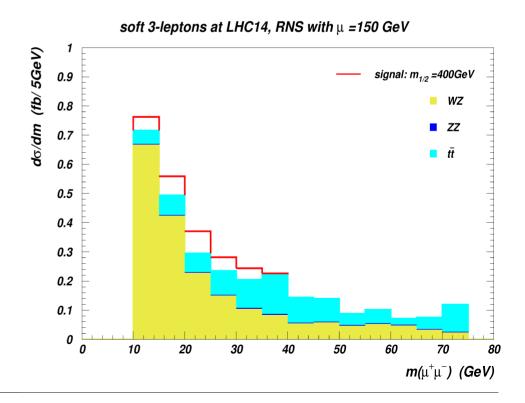
- Mass gap between higgsinos is <30 GeV \Rightarrow very soft leptons for higgsino pair production. But ATLAS/CMS can detect soft leptons $p_T(\mu) > 5~{
 m GeV},~p_T(e) > 10~{
 m GeV}$
- $pp \longrightarrow \widetilde{W}_1 \widetilde{Z}_2 \longrightarrow (e\nu_e \widetilde{Z}_1) + (\mu^+ \mu^- \widetilde{Z}_1)$
- For $m_{1/2} < 400-500~{
 m GeV}$ and $\mu{=}150{
 m GeV}$ most e pass trigger threshold



Soft trileptons

arXiv: 1302.5816

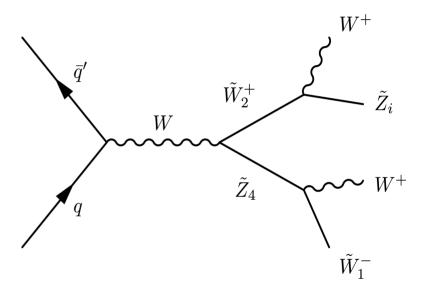
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- $pp \longrightarrow \widetilde{W}_1 \widetilde{Z}_2 \longrightarrow (e\nu_e \widetilde{Z}_1) + (\mu^+ \mu^- \widetilde{Z}_1)$
- For $m_{1/2} < 400-500~{
 m GeV}$ and $\mu{=}150{
 m GeV}$ most e pass trigger threshold
- A shape analysis may allow to claim a signal with high luminosity - confirmation channel



Same-sign Dibosons

arXiv: 1302.5816

- Sizeable production cross section (~10-100 fb) for wino-like $\widetilde{Z}_4\widetilde{W}_2$ and $\widetilde{W}_2\widetilde{W}_2$
- Each decays into W about 50%
- Two W's are not charge correlated
- Same-sign WW is novel signature, charcteristic of light higgsinos
- No 2->2 SM production of ssWW
- BG from WZ and ttbar removed by mT>125 GeV cut
- 2155+MET is not probed by existing di-lepton searches that require significant hadronic activity

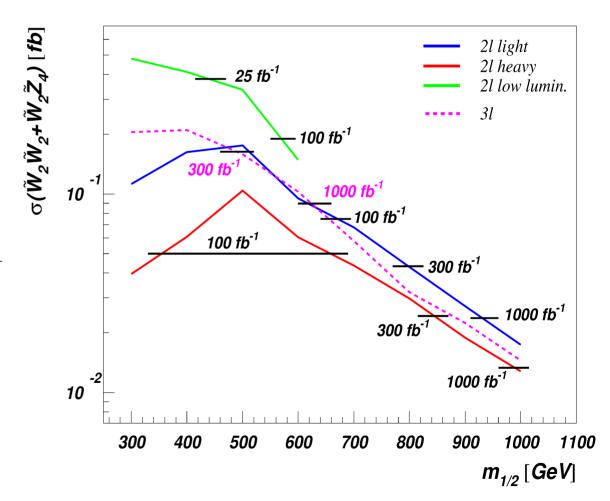


ssWW prospects

arXiv: 1302.5816, 1310.4858

- Reach for 100 (1000) fb⁻¹ extends to $m_{1/2} \sim 680~(1000)~{
 m GeV}$
- Much better than canonical 3-lepton reach
- Exceeds direct gluino production reach, $m_{1/2} \sim 700 {
 m GeV~with}~300 {
 m fb}^{-1}$ (if gaugino masses are unified)
- Independent of gluino search!

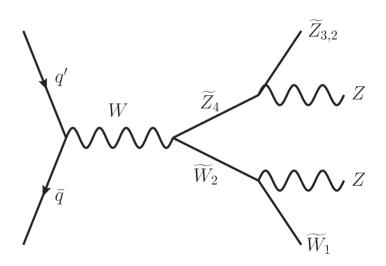
NUHM2: m_0 =5 TeV, A_0 =-1.6 m_0 , $tan\beta$ =15, μ =150 GeV, m_A =1 TeV



4-lepton signature

arXiv: 1310.4858

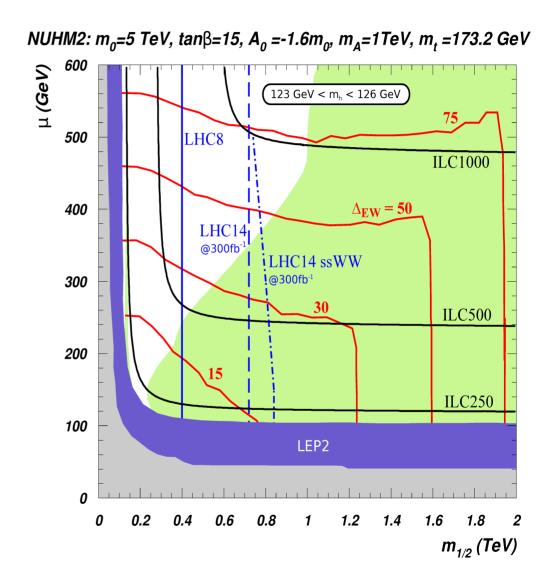
- Winos also have significant BF to Z, 25-50%
- Novel channel for low $|\mu|$ models: 4 isolated leptons + MET and no jets
- BGs controlled by MET > 200 GeV cut
- Reach for 300 (1000) fb⁻¹ extends to $m_{1/2} \sim 500~(650)~{\rm GeV}$ comparable to canonical 3-leptons -> confirmatory channel for same-sign WW
- Current 4-lepton searches are optimized for glunio cascades in RPV with many leptons and jets



LHC-ILC complimentarity

arXiv: 1310.4858

- In light higgsino models, ssWW offers the largest reach $m_{1/2} \sim 1~{
 m TeV}$ for HL-LHC
- Still not enough to fully probe
 <3% EWFT region
- ILC with $\sqrt{s}\sim 600~{\rm GeV}$ is needed for complete check of FT in MSSM
- ILC can measure higgsino masses with sub-GeV precision (see arXiv:1404.7510)

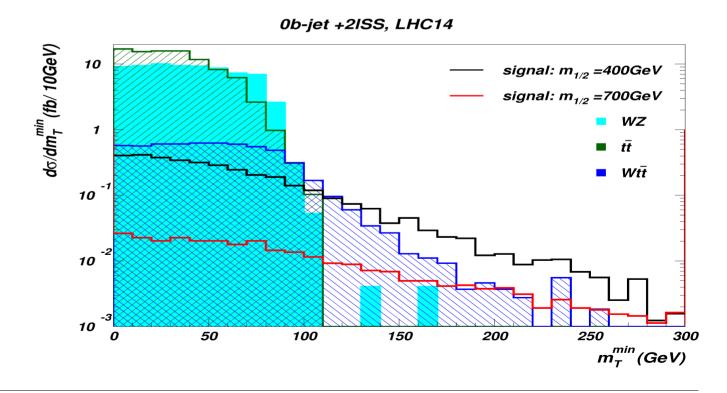


Conclusions

- Although FT is intrinsically model dependent, generic pheno studies of natural SUSY are possible with $\Delta_{_{\text{FW}}}$ lower bound on FT.
- Small μ is necessary (but not sufficient) condition for naturalness. It is more fundamental than light stops!
- Higgsino-like chargino and neutralinos are difficult to detect at LHC due to low visible energy release from their decays.
- Diboson productions are novel signatures characteristic of light higgsino scenario. They depend only on EW-ino spectrum.
- Same-sign WW can probe wino masses up to 550 (800) GeV, for 100 (1000) fb⁻¹ at LHC14. Larger reach than in canonical trilepton channel.
- LHC can cover large portion of para space by ssWW, but ILC with ~600 GeV is necessary to complete the test of SUSY naturalness with FT >3%.

Cuts for ssWW

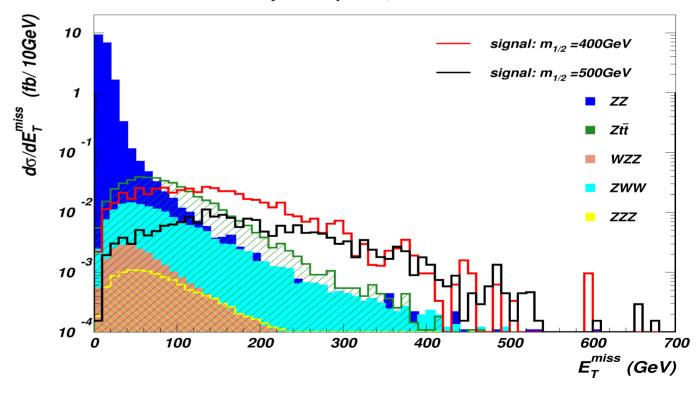
- b-jet veto (60% eff.)
- ullet 2 isolated same-sign leptons $p_T(l_1)>20GeV,\; p_T(l_2)>10GeV$
- $m_T^{min} \equiv \min\left[m_T(l_1, E_T), m_T(l_2, E_T)\right] > 125 GeV$ removes WZ and ttbar due to kinematic cutoff for on-shell W
- MET>200GeV



Cuts for ZZ

- b-jet veto (60% eff.)
- 4 isolated leptons with $p_T(l)>10~{
 m GeV}, \ |\eta(l)|<2.5$
- MET > 100 (200) GeV





Higgsinos mass gap

arXiv: 1404.7510

